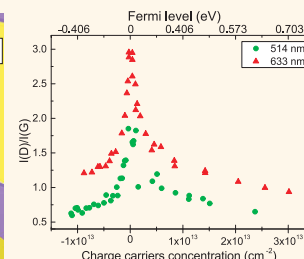
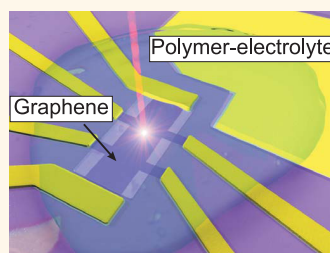


# Doping Dependence of the Raman Spectrum of Defected Graphene

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**ABSTRACT** We investigate the evolution of the Raman spectrum of defected graphene as a function of doping. Polymer electrolyte gating allows us to move the Fermi level up to 0.7 eV, as directly monitored by *in situ* Hall-effect measurements. For a given number of defects, we find that the intensities of the D and D' peaks decrease with increasing doping. We assign this to an increased total scattering rate of the photoexcited electrons and holes, due to the doping-dependent strength of electron–electron scattering. We present a general relation between D peak intensity and defects valid for any doping level.



**KEYWORDS:** graphene · Raman spectroscopy · defects

Raman spectroscopy is one of the most used characterization techniques in carbon science and technology.<sup>1</sup> The measurement of the Raman spectrum of graphene<sup>2</sup> triggered a huge effort to understand phonons, electron–phonon ( $e-ph$ ), magnetophonon and electron–electron ( $e-e$ ) interactions in graphene, as well as the influence of the number and orientation of layers, electric or magnetic fields, strain, doping, disorder, quality and types of edges, and functional groups.<sup>3</sup>

Quantifying defects in graphene is crucial both to gain insight in fundamental properties, and for applications. Ref 4 introduced a three-stage classification of disorder, leading from graphite to amorphous carbons, that allows to simply assess all the Raman spectra of carbons. Stage 1: graphene to nanocrystalline graphene. Stage 2: nanocrystalline graphene to low- $sp^3$  amorphous carbon. Stage 3: low- $sp^3$  amorphous carbon to high- $sp^3$  amorphous carbon. Here we focus on stage 1, the most relevant when considering the vast majority of publications dealing with graphene production, processing and applications. In stage 1 the Raman spectrum evolves as follows:<sup>4</sup> (a) the D peak appears and the ratio of the D and G peak intensities,  $I(D)/I(G)$ , increases; (b) the D' peak appears; (c) all peaks broaden; (d) the D + D' peak appears; (e) at the end of stage 1, the G and D' peaks are so wide that it is sometimes more

convenient to consider them as a single, up-shifted, wide G band at  $\sim 1600\text{ cm}^{-1}$ .

In their seminal work, Tuinstra and Koenig noted that  $I(D)/I(G)$  varied inversely with the crystal size,  $L_a$ :  $I(D)/I(G) = C(\lambda)/L_a$ , where  $C(514\text{ nm}) \sim 4.4\text{ nm}^{5-7}$  ( $\lambda$  being the excitation wavelength). Initially, this was interpreted in terms of phonon confinement: the intensity of the forbidden process would be ruled by the “amount of lifting” of the selection rule,<sup>5</sup>  $\Delta q \propto 1/\Delta x$ , with  $\Delta x \approx L_a$ . Now, it is understood theoretically and established experimentally, that the D peak is produced only in a small region of the crystal (size  $\sim v_F/\omega_D \sim 3-4\text{ nm}$ , where  $v_F$  is the Fermi velocity and  $\omega_D$  is the phonon frequency) near a defect or an edge.<sup>8-10</sup> For a nanocrystallite,  $I(G)$  is proportional to the sample area,  $\propto L_a^2$ , while  $I(D)$  is proportional to the overall length of the edge, which scales as  $\propto L_a$ . Thus,  $I(D)/I(G) \propto 1/L_a$ . For a sample with rare defects,  $I(D)$  is proportional to the total number of defects probed by the laser spot. Thus, for an average interdefect distance  $L_D$ , and laser spot size  $L_L$ , there are on average  $(L_L/L_D)^2$  defects in the area probed by the laser, then  $I(D) \propto (L_L/L_D)^2$ . On the other hand,  $I(G)$  is proportional to the total area probed by the laser  $\propto (L_L)^2$ , thus  $I(D)/I(G) = C''(\lambda)/L_D^2$ . For very small  $L_D$ , one must have  $C''(\lambda)/L_D^2 = I(D)/I(G) = C(\lambda)/L_a$ . This condition gives an estimate of  $C''(514\text{ nm}) \sim 90\text{ nm}$ . Ref 9 measured  $I(D)/I(G)$  for irradiated single layer

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graphene (SLG) with known  $L_D$ , derived from scanning tunneling microscopy, obtaining  $I(D)/I(G) \approx 145/L_D^2$  at 514 nm excitation, in agreement with this simple estimate.

Ref 11 then considered the excitation energy dependence of the peaks areas and intensities, for visible excitation energy. A fit to the experimental data gave the relation<sup>11</sup>

$$L_D^2[\text{nm}^2] = \frac{4.3 \times 10^3}{E_L^4[\text{eV}^4]} \left[ \frac{I(D)}{I(G)} \right]^{-1} \quad (1)$$

where  $E_L$  is the laser excitation energy. By considering point-like defects, separated from each other by  $L_D[\text{nm}]$ , eq 1 can be restated in terms of defect density  $n_D[\text{cm}^{-2}] = 10^{14}/\{\pi L_D^2[\text{nm}^2]\}$ :<sup>11</sup>

$$n_D[\text{cm}^{-2}] = 7.3 \times 10^9 E_L^4[\text{eV}^4] \frac{I(D)}{I(G)} \quad (2)$$

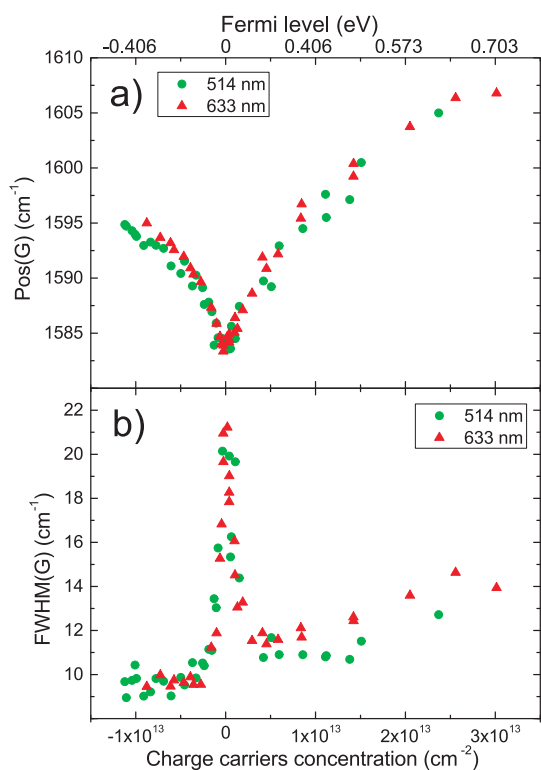
Note that these relations are limited to Raman-active defects. Perfect zigzag edges,<sup>8,12</sup> charged impurities,<sup>13,14</sup> intercalants,<sup>15</sup> uniaxial and biaxial strain<sup>16,17</sup> do not generate a D peak. For these types of “silent” defects, other Raman signatures can be used. A perfect edge does change the G peak shape,<sup>18,19</sup> while strain, intercalants, and charged impurities have a strong influence on the G and 2D peaks.<sup>13–16</sup> In this case, the combination of Raman spectroscopy with other independent probes of the number of defects can provide a wealth of information on the nature of such defects.

We note as well that these relations are derived assuming negligible Fermi level,  $E_F$ , shift. It is known that doping has major effects on the Raman spectra.<sup>14,20,21</sup> The G peak position,  $\text{Pos}(G)$ , increases and its Full Width at Half Maximum,  $\text{FWHM}(G)$ , decreases for both electron ( $e$ ) and hole ( $h$ ) doping. The G peak stiffening is due to the nonadiabatic removal of the Kohn anomaly at the Brillouin Zone (BZ) center,  $\Gamma$ .<sup>20</sup> The  $\text{FWHM}(G)$  sharpening is due to Pauli blocking of phonon decay into  $e-h$  pairs when the  $e-h$  gap is higher than the phonon energy,<sup>20,22</sup> and saturates for  $E_F$  bigger than half phonon energy.<sup>20,22</sup> Furthermore, in SLG the ratio of the heights of the 2D and G peaks,  $I(2D)/I(G)$ , and their areas,  $A(2D)/A(G)$ , is maximum for zero doping,<sup>2,23</sup> and decreases for increasing doping. The doping dependence of the 2D intensity results from its sensitivity to the scattering of the photoexcited  $e$  and  $h$ . Assuming the dominant sources of scattering to be phonon emission and  $e-e$  collisions, ref 21 showed that, while the former is not sensitive to doping, the latter is. Then, the doping dependence of the 2D peak can be used to estimate the corresponding  $e-ph$  coupling.<sup>21</sup> These considerations apply for  $|E_F|$  small compared to  $\hbar\omega_L/2$  ( $\omega_L$  being the angular frequency of the incident photon). In the past few years, much higher doping levels have been achieved.<sup>15,24,25</sup> One of the effects of high doping is the increase in  $I(G)$ . Doping changes the occupations of electronic states

and, since transitions from an empty state or to a filled state are impossible, it can exclude some BZ regions from contributing to the Raman matrix element. Because of suppression of destructive interference, this leads to an enhancement of the G peak when  $|E_F|$  matches  $\hbar\omega_L/2$ , as predicted theoretically<sup>26</sup> and observed experimentally.<sup>24,25</sup> Another effect of high doping is on the 2D peak, which is suppressed when the conduction band becomes filled at the energy probed by the laser.<sup>15</sup> There are three cases: (i) when  $\omega_L\omega_{sc} > 2|E_F|/\hbar$  (where  $\omega_{sc}$  is the angular frequency of the emitted photon), all processes are allowed and the 2D band is observed, (ii) when  $\omega_{sc} < 2|E_F|/\hbar < \omega_L$ , the photon absorption is allowed but the phonon emission is excluded by Pauli blocking; (iii) when  $\omega_L\omega_{sc} < 2|E_F|/\hbar$ , both photon absorption and phonon emission are blocked. Therefore, only when  $2|E_F|/\hbar < [\omega_L - \text{Pos}(2D)]$ , the 2D band is observable.

While a significant effort was devoted to understanding the effect of defects in samples with negligible doping,<sup>4,8,11</sup> and the effect of doping in samples with negligible defects,<sup>14,15,21,24–26</sup> the combined effect of doping and defects on the Raman spectrum of SLG has received little attention. However, most samples produced by either micromechanical exfoliation, chemical vapor deposition, liquid phase exfoliation or carbon segregation from SiC or metal substrates are naturally doped due to the extreme sensitivity of graphene to the presence of adsorbates (*e.g.*, moisture) and to the interaction with the underlying substrate.<sup>13,14,27</sup> Additionally, many of them also have defects, or defects may appear during processing for device integration. It is thus critical to understand if and how defects can be detected and quantified by Raman spectroscopy in doped samples.

Here we study the dependence of the Raman spectrum of defected SLG on the level of electrostatic doping, in samples with a fixed amount of defects. We combine polymer electrolyte gating<sup>14</sup> with *in situ* Hall-effect measurements and Raman spectroscopy at different excitation wavelengths. This allows us to vary  $E_F$  from  $\approx -0.7$  eV ( $h$ -doping) up to  $\approx 0.4$  eV ( $e$ -doping), a much wider span than what can be achieved by the common 300 nm  $\text{SiO}_2$  back gate (usually restricted to  $\approx \pm 0.3$  eV because of the limited gate capacitance<sup>28</sup>) and large enough to cover the range of doping found in the vast majority of papers in literature. This range is however smaller than  $\hbar\omega_L/2$ , to exclude additional effects due to Pauli blocking on the peaks' intensities, again consistent with the case in most papers. We find that the intensity of defect-related peaks, D and D', strongly decreases with  $E_F$ . We assign this to increased broadening of the electronic states due to increased  $e-e$  scattering for higher doping. We then modify eqs 1 and 2 and give general relations between D peak intensity and defects valid for any doping level.



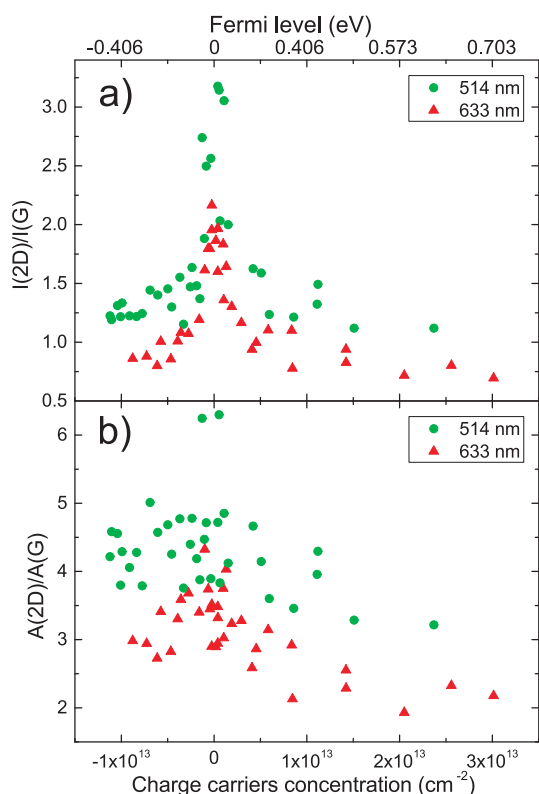
**Figure 1.** Dependence of (a) Pos(G) and (b) FWHM(G) on  $E_F$  and carrier concentration for excitation at 514 and 633 nm.

## RESULTS AND DISCUSSION

The sample preparation and the details of the polymer electrolyte gating, *in situ* Hall-effect and Raman spectroscopy measurements at different excitation wavelengths are discussed in Methods.

Figure 1 plots Pos(G) and FWHM(G) as a function of  $E_F$ . As  $E_F$  moves from the charge neutrality point, the G peak blue-shifts and narrows, consistent with what reported in pristine graphene in the presence of moderate electrostatic doping<sup>14,20</sup> with  $E_F \ll \hbar\omega_L/2$ . The hardening of the G mode is due to nonadiabatic removal of a Kohn anomaly at  $\Gamma$ ,<sup>20,29</sup> and the reduction of width is due to Pauli exclusion principle inhibiting phonon decay into  $e-h$  pairs when  $E_F$  surpasses half the phonon energy.<sup>20</sup> Figure 2 shows the doping dependence of the intensity and area ratio of 2D and G. Both decrease with increasing  $E_F$  due to the effect of increased  $e-e$  interaction.<sup>21</sup>  $A(G)$  remains roughly constant with  $E_F$ ,<sup>21</sup> while  $I(G)$  is reduced close to the charge neutrality point due to the damping of the phonon decaying into  $e-h$  pairs which increases FWHM(G). This results in a stronger doping dependence of  $I(2D)/I(G)$  than  $A(2D)/A(G)$ , as for Figure 2.

Figure 3 plots the dependence of Pos(2D) on  $E_F$ . For  $h$ -doping, Pos(2D) increases, while for  $e$ -doping, Pos(2D) slightly increases at first, then decreases as  $E_F$  keeps rising.<sup>14</sup> This is due to a modification of the lattice parameters caused by doping, which changes the total number of charges, with a consequent stiffening/softening of

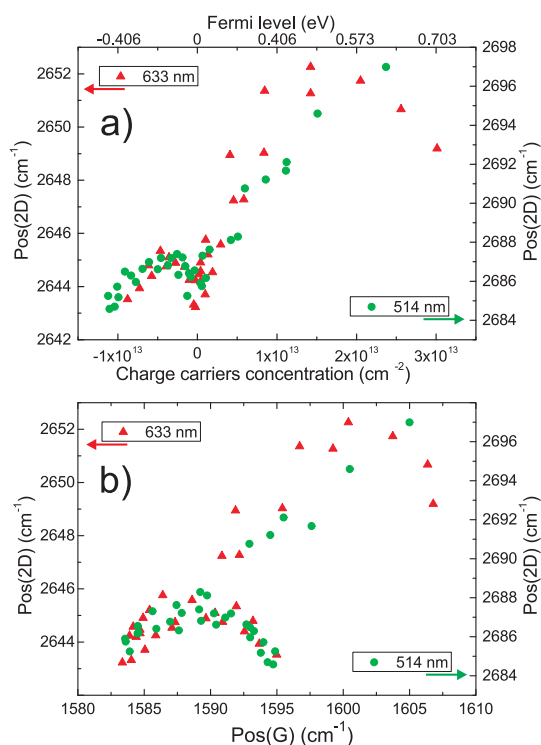


**Figure 2.** Dependence of (a)  $I(2D)/I(G)$  and (b)  $A(2D)/A(G)$  on  $E_F$  and carrier concentration for excitation at 514 and 633 nm.

the phonons.<sup>14</sup> The dependence of Pos(2D) on  $E_F$  is different from that of Pos(G). Indeed, the latter always increases with  $E_F$ , as highlighted in Figure 3b, where Pos(2D) is plotted against Pos(G). This allows to distinguish  $e^-$  from  $h^-$  doping in graphene using Raman spectroscopy.

Now, we focus on the doping dependence of the main Raman signatures of defects in graphene: D and D'. We consider two samples, A, B. The defect concentration in is evaluated from  $I(D)/I(G)$  measured at low doping ( $n < 5 \times 10^{11} \text{ cm}^{-2}$ ,  $E_F < 90 \text{ meV}$ ), with eq 2. This gives  $n_D^A \sim 4.4 \times 10^{11} \text{ cm}^{-2}$  and  $n_D^B \sim 2.5 \times 10^{11} \text{ cm}^{-2}$ . Figure 4 plots the Raman spectra of sample A as a function of doping. Figure 5a,b shows that the evolution of the D peak with doping is similar to that of the 2D peak [compare Figure 5a,b with Figure 2a,b], with a marked decrease in  $I(D)/I(G)$ , for increasing  $E_F$ , Figure 5a. The evolution of the peak's area ratio [Figure 5b], more robust with respect to various perturbations of the phonon states than the height,<sup>21</sup> shows a decrease with  $E_F$ . The same behavior is observed for  $I(D')/I(G)$  and  $A(D')/A(G)$ , Figure 5c,d. Figure 5e,f also indicates that  $I(D)/I(2D)$  and  $A(D)/A(2D)$  have no clear dependence on  $E_F$ , pointing to a similarity of the physical phenomena determining the doping-dependent Raman scattering for these two different Raman processes.

Similar trends as in Figure 5 are also observed for the less defective sample B.



**Figure 3.** (a) Dependence of Pos(2D) on  $E_F$  and carrier concentration and (b) Pos(2D) as a function of Pos(G) for excitation at (left y-axis) 633 nm and (right y-axis) 514 nm.

The decrease in  $I(D)/I(G)$  with  $E_F$  can be described by a power relation. Rescaling  $I(D)/I(G)$  by the fourth power of  $E_L$  and normalizing by the amount of defects as measured at low doping, all the values collapse on the same line with slope  $\alpha \sim -0.54$  in a logarithmic plot against the absolute value of  $E_F$  (see Figure 6).

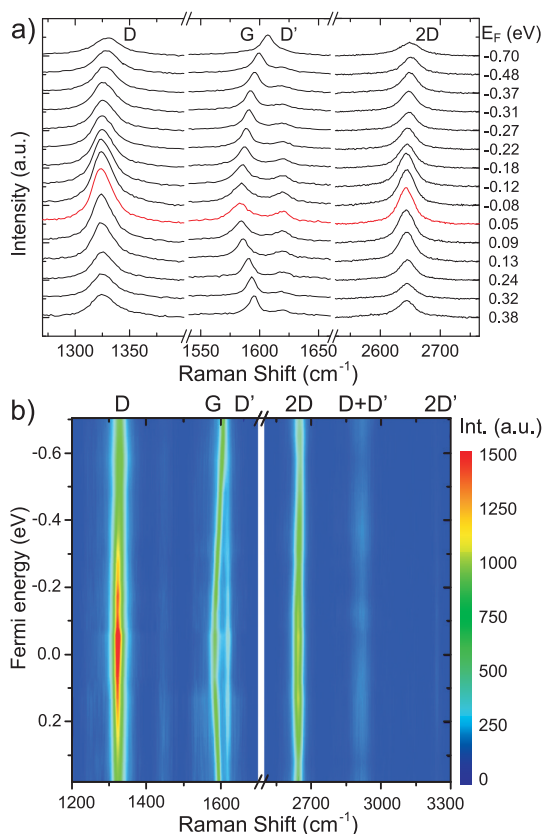
We can thus modify eq 1 for samples with non-negligible doping:

$$L_D^2[\text{nm}^2] = \frac{(1.2 \pm 0.3) \times 10^3}{E_L^4[\text{eV}^4]} \left[ \frac{I(D)}{I(G)} \right]^{-1} \{E_F[\text{eV}]\}^{-0.54 \pm 0.04} \quad (3)$$

and eq 2:

$$n_D[\text{cm}^{-2}] = (2.7 \pm 0.8) \times 10^{10} E_L^4[\text{eV}] \frac{I(D)}{I(G)} \{E_F[\text{eV}]\}^{0.54 \pm 0.04} \quad (4)$$

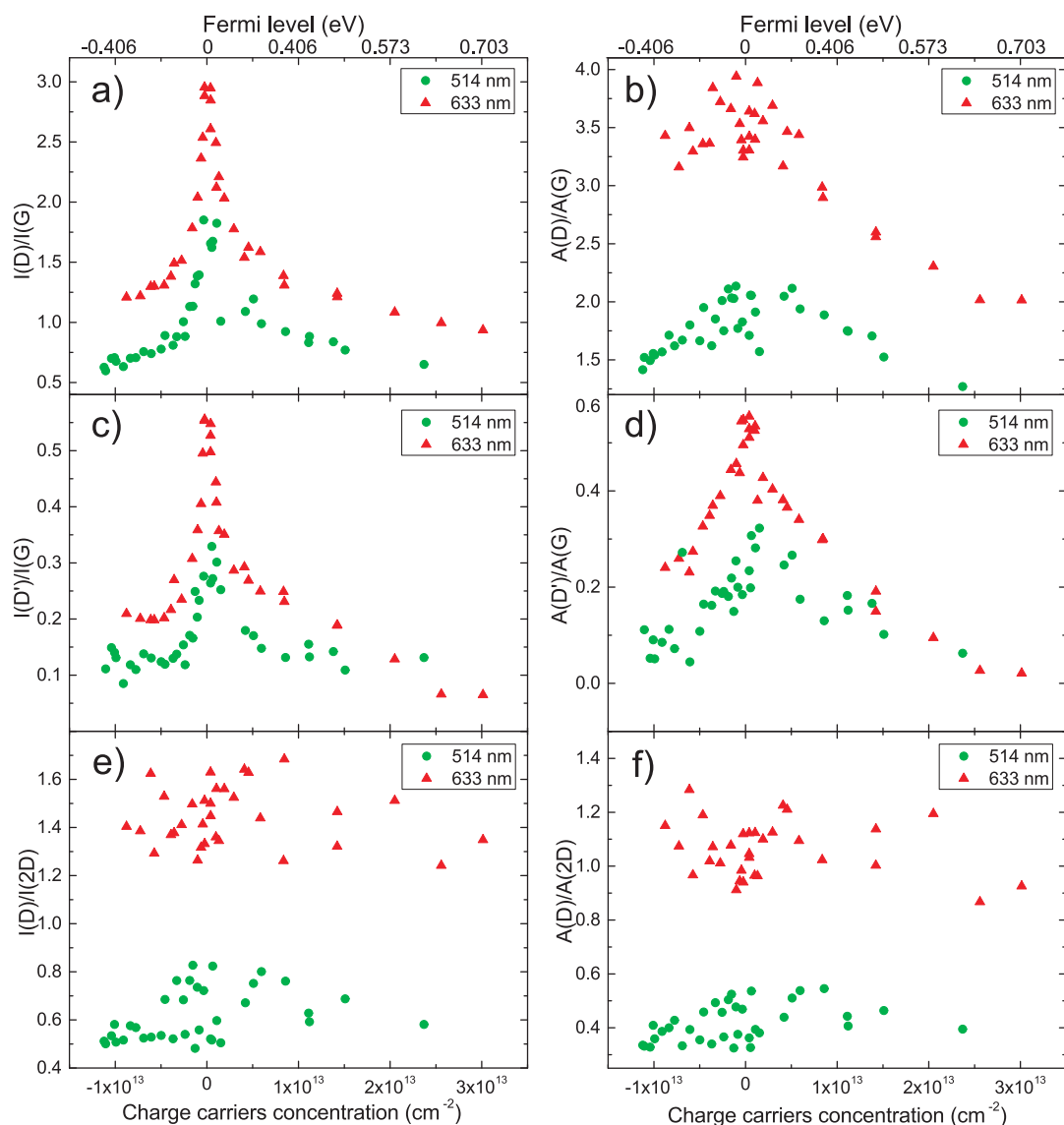
Equations 3 and 4 are valid for samples with a defect concentration corresponding to stage 1, by far the most relevant for graphene production and applications, and for  $E_F < E_L/2$ , in order to avoid Pauli blocking effects on the intensity of peaks.<sup>15,24–26</sup> Combining  $I(D)/I(G)$  and FWHM(G) it is possible to discriminate between stages 1 or 2, because samples in stage 1 and 2 could have the same  $I(D)/I(G)$ , but not the same FWHM(G), which is much larger in stage 2.<sup>4,11</sup> Since most graphene samples in literature show doping levels  $\sim 200$ – $500$  meV, and Raman characterization is mostly carried out with excitation wavelengths in the visible ( $1 \text{ eV} > \hbar\omega_L/2 > 1.5 \text{ eV}$ ),



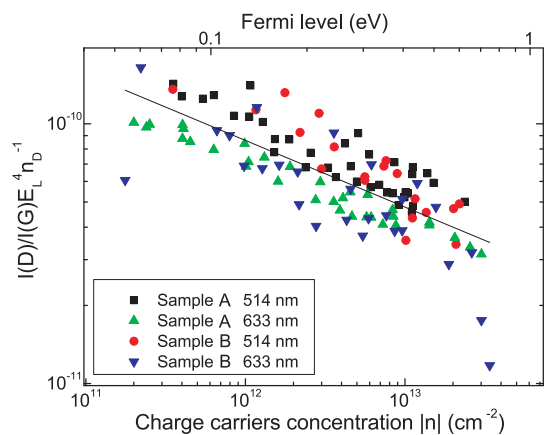
**Figure 4.** (a) Raman spectrum of defected graphene, measured at 633 nm, for different  $E_F$ . The spectra are normalized with respect to  $I(G)$ . The spectrum highlighted in red corresponds to the lowest  $E_F$  probed in our experiments, separating  $h$ -doping (negative  $E_F$ ) and  $e$ -doping (positive  $E_F$ ). (b) Contour plot of the same data, showing the intensity of the Raman signal as a function of  $E_F$  and Raman shift.

eqs 3 and 4 cover the vast majority of experimental conditions in graphene science and technology.<sup>27</sup>

It is useful to consider a practical example on how eqs 3 and 4 may be used. Figure 7b plots a typical Raman spectrum of a defected doped sample compared to the spectrum of the same sample, therefore with the same amount of defects, in the undoped case (Figure 7a). In absence of external means to derive  $E_F$ , one could use Raman spectroscopy to evaluate  $E_F$ . Since Pos(G), FWHM(G), Pos(2D) and  $I(2D)/I(G)$  for defected graphene within stage 1 evolve consistently with what reported for non defective samples,<sup>14</sup> one can use these to estimate  $E_F \leq 100$  meV in case (a) and  $E_F \sim 500$  meV in case (b). Comparing Pos(2D) with Pos(G) and using Figure 3b, it is possible to conclude that the sample is  $h$ -doped. Since  $I(D)/I(G) \sim 1.24$  at 633 nm, if we ignore doping and use eqs 2 and 1 we get  $n_D \approx 1.3 \times 10^{11} \text{ cm}^{-2}$  or  $L_D \approx 16$  nm. Equations 4 and 3 instead give  $n_D \approx 3.4 \times 10^{11} \text{ cm}^{-2}$  or  $L_D \approx 10$  nm. The defect density estimated taking into account doping is more than twice that from eq 2. This can make a difference for the optimization of methods of production and processing of graphene, especially for what concerns particular applications, such as transparent



**Figure 5.** (a)  $I(D)/I(G)$ , (b)  $A(D)/A(G)$ , (c)  $I(D')/I(G)$ , (d)  $A(D')/A(G)$ , (e)  $I(D)/I(2D)$  (f)  $A(D)/A(2D)$  as a function of  $E_F$  (top x-axis) or charge carrier concentration (bottom x-axis) at two different excitation wavelengths for sample A, with  $n_D^A \sim 4.4 \times 10^{11} \text{ cm}^{-2}$ .



**Figure 6.** Doping dependence of  $I(D)/I(G)$  for both samples A and B rescaled by the fourth power of  $E_L$  and normalized by the amount of defects.

conductive films, where low sheet resistance may be achieved through high doping.<sup>30</sup>

We note that ref 31 reported a doping dependence of  $I(D)$ , which was found to sharply increase for  $|E_F| \sim \hbar\omega_L/2$ . This was assigned to quantum interference between different Raman pathways, similarly to the case of the G peak in refs 24–26. However, ref 31 suffers from poor accuracy in the determination of the charge carriers concentration, leading to results inconsistent with literature.<sup>14,25</sup> Indeed, in ref 31 the graphene Raman peaks are found to evolve with doping as in ref 31, but at doping levels almost one order of magnitude higher; *e.g.*,  $\text{Pos}(2D) \approx 2670 \text{ cm}^{-1}$  for  $E_L = 2.41 \text{ eV}$  is reached in ref 14 for  $n \approx 3.4 \times 10^{13} \text{ cm}^{-2}$  while in ref 31 is observed for  $n \approx 2.6 \times 10^{14} \text{ cm}^{-2}$ . The inconsistency of ref 31 is evident also in the analysis of



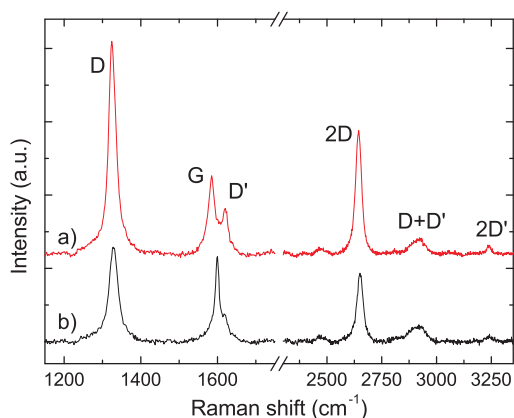


Figure 7. Raman spectra of sample A, acquired at 633 nm, for (a)  $E_F \leq 100$  meV and (b)  $E_F \sim 500$  meV.

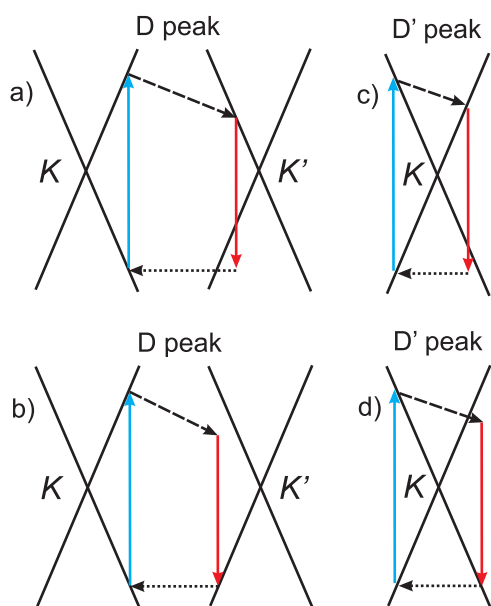


Figure 8. Raman processes giving rise to the D and D' peaks.<sup>3</sup> (a,b) Defect-assisted intervalley one-phonon D peak. (c,d) Defect-assisted intravalley one-phonon D' peak. Vertical solid arrows represent interband transitions accompanied by photon absorption (upward arrows) or emission (downward arrows). Dashed arrows and horizontal dotted arrows represent phonon emission and elastic defect scattering, respectively.

the doping dependence of  $\text{Pos}(G)$ . Ref 14 reported for pristine graphene  $\delta\text{Pos}(G)/\delta E_F \approx 30 \text{ cm}^{-1} \text{ eV}^{-1}$  for  $e$ -doping and  $\delta\text{Pos}(G)/\delta E_F \approx 42 \text{ cm}^{-1} \text{ eV}^{-1}$  for  $h$ -doping, as also confirmed by ref 25 and our present work, where we obtain  $\delta\text{Pos}(G)/\delta E_F \approx 30 \text{ cm}^{-1} \text{ eV}^{-1}$  for  $e$ -doping and  $\delta\text{Pos}(G)/\delta E_F \approx 39 \text{ cm}^{-1} \text{ eV}^{-1}$  for  $h$ -doping (see Figure 1). Ref 31 has a much weaker doping dependence of  $\text{Pos}(G)$ , in particular  $\delta\text{Pos}(G)/\delta E_F \approx 12 \text{ cm}^{-1} \text{ eV}^{-1}$  for  $e$ -doping and  $\delta\text{Pos}(G)/\delta E_F \approx 19 \text{ cm}^{-1} \text{ eV}^{-1}$  for  $h$ -doping. This indicates that the actual doping level reached by ref 31 is much smaller than what claimed and surely far from  $|E_F| \approx \hbar\omega_L/2 \geq 1.2 \text{ eV}$  necessary to achieve the blocking of Raman pathways<sup>24–26</sup> for  $\lambda = 514.5 \text{ nm}$  (as used in ref 31), thus

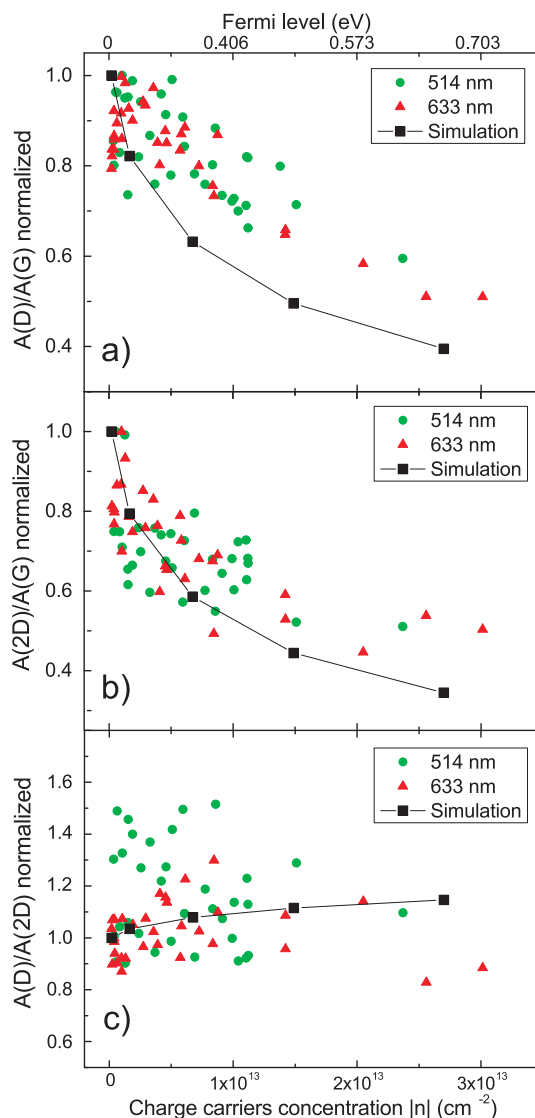


Figure 9. Evolution of (a)  $A(D)/A(G)$ , (b)  $A(2D)/A(G)$ , (c)  $A(D)/A(2D)$  as a function of doping for sample A, normalized with respect to their values at low doping, compared to the calculated values. Green circles and red triangles show the experimental data at two different excitation wavelengths and the black squares indicate the calculated values. The line is to guide the eye.

compromising the basis of their physical explanation for the observed increase of  $I(D)$ .

Ref 32 reported the  $I(D)$  dependence on back-gate bias in SLG, observing an increase in  $I(D)/I(G)$  with increasing gate voltage, in principle the opposite of what we report here in Figure 5a. Ref 32 attributed this to a change in the total amount of defects in graphene due to electrochemical reactions involving the water layer trapped at the interface between graphene and the silicon dioxide substrate. This is quite a different case with respect to that studied here, where the number of defects is kept constant as a function of doping, as confirmed by the repeatability of the Raman measurements through several gate voltage sweeps.

In order to understand the physical reason for the D peak decrease with  $E_F$ , we consider the Raman scattering processes in more detail. The G peak corresponds to the high frequency  $E_{2g}$  phonon at  $\Gamma$ . The D peak is due to the breathing modes of six-atom rings and requires a defect for its activation.<sup>4,5,33</sup> It comes from TO phonons around the BZ edge  $\mathbf{K}$ ,<sup>4,5</sup> is active by double resonance (DR),<sup>33</sup> and is strongly dispersive with excitation energy,<sup>34</sup> due to a Kohn Anomaly at  $\mathbf{K}$ .<sup>35</sup> DR can also happen as intravalley process, *i.e.*, connecting two points belonging to the same cone around  $\mathbf{K}$  (or  $\mathbf{K}'$ ). This gives the D' peak. The dominant scattering paths in the D and D' processes are illustrated in Figure 8. The 2D peak is the D peak overtone. The 2D' peak is the D' overtone. Since 2D and 2D' originate from a process where momentum conservation is satisfied by two phonons with opposite wavevectors, no defects are required for their activation, and are thus always present.<sup>3,21</sup>

The Raman spectra can be modeled within second-order perturbation theory, by summing over all scattering pathways, expressed with the Fermi's golden rule to the fourth order.<sup>36</sup> Each of these pathways consists of the creation of an  $e-h$  pair due to the interaction with incident excitation laser, one  $e-ph$  and one electron-defect ( $e-def$ ) scattering, and the  $e-h$  pair recombination.<sup>3,37</sup> The amplitude of each pathway is given by the corresponding matrix element. By summing over all possible processes in the BZ, constructive and destructive interferences between the different quantum paths are fully taken into account. Both the photoexcited  $e$  and  $h$  may scatter, resulting in four different combinations,  $ee$ ,  $eh$ ,  $he$ , and  $hh$ . Taking into account the two relative orderings for the  $e-ph$  and  $e-def$  scattering, there are eight different contributions for each  $\mathbf{k}$ -point in the BZ; *e.g.*, the matrix element for a phonon-defect scattering pathway with the electron first scattering and emitting a phonon, and the hole then scattering with a defect, is given by<sup>37</sup>

$$K_{\text{ehl}}^{\text{pd}} = \frac{\langle \mathbf{k} + \mathbf{q}, \pi | H_{e-em, \text{out}} | \mathbf{k} + \mathbf{q}, \pi^* \rangle \langle \mathbf{k}\pi | H_D | \mathbf{k} + \mathbf{q}, \pi \rangle \langle \mathbf{k}, \pi^* | \Delta H_{\mathbf{q}\nu} | \mathbf{k}\pi^* \rangle \langle \mathbf{k}\pi^* | H_{e-em, \text{in}} | \mathbf{k}\pi \rangle}{(E_L - E_{\mathbf{k}+\mathbf{q}}^{\pi^*} + E_{\mathbf{k}+\mathbf{q}}^{\pi} - \hbar\omega_{-\mathbf{q}}^{\nu} - 2i\gamma_{\text{tot}})(E_L - E_{\mathbf{k}+\mathbf{q}}^{\pi^*} + E_{\mathbf{k}}^{\pi} - \hbar\omega_{-\mathbf{q}}^{\nu} - 2i\gamma_{\text{tot}})(E_L - E_{\mathbf{k}}^{\pi^*} + E_{\mathbf{k}}^{\pi} - 2i\gamma_{\text{tot}})} \quad (5)$$

where  $E_{\mathbf{k}}^{\pi}$  is the energy of the electronic state in the valence (conduction) band at wave vector  $\mathbf{k}$ ,  $\omega_{\mathbf{q}}^{\nu}$  is the phonon frequency of branch  $\nu$  at phonon momentum  $\mathbf{q}$ , and  $\gamma_{\text{tot}}$  describes the broadening of the intermediate electronic states. The nominator contains matrix elements corresponding to the light absorption and emission ( $H_{e-em, \text{in}}$  and  $H_{e-em, \text{out}}$ , respectively),  $e-ph$  scattering ( $\Delta H_{\mathbf{q}\nu}$ ) and defect scattering ( $\Delta H_D$ ). In the case of the 2D peak, the defect scattering matrix element is replaced by a second  $e-ph$  scattering.<sup>3,37</sup>

The Raman intensity is then obtained as

$$I_D(\omega) \propto \sum_{\mathbf{q}, \nu} \left| \sum_{\mathbf{k}, \alpha} K_{\alpha}^{\text{pd}}(\mathbf{k}, \mathbf{q}, \nu) \right|^2 \delta(\omega - \omega_{\mathbf{q}}^{\nu}) \quad (6)$$

and similarly for the 2D peak, but involving a summation over two phonon branches  $\nu$  and  $\mu$ .

The contributions to the total broadening  $\gamma_{\text{tot}}$  due to  $e-ph$ ,  $\gamma_{e-ph}$ , and  $e-def$  scattering,  $\gamma_D$ , are determined by the corresponding scattering matrix elements (see ref 37). In the presence of doping, a contribution due to  $e-e$  scattering,  $\gamma_{ee}$  increasing with  $E_F$ , has to be included.<sup>21</sup> In principle, these quantities depend on the wave vector and band index of the electronic state; thus, they are inhomogeneous in the BZ. For simplicity, here, as in ref 37, they are assumed to be independent of the electronic state.

Figure 5e,f indicates a similar doping dependence of the D and 2D peaks. The  $A(2D)$  decrease with increasing doping is due to additional broadening of the intermediate  $e-h$  states induced by  $e-e$  interactions, with  $\gamma_{ee} \propto |E_F|$ .<sup>21</sup> Thus, the total broadening at a given  $E_F$  is given by  $\gamma_{\text{tot}} = \gamma_{e-ph} + \gamma_D + \gamma_{ee}(E_F)$ , where  $\gamma_{ee}(E_F) = 0.06|E_F|$ .<sup>21</sup> For the numerical simulation using eq 6, the band structure,  $e-ph$  matrix elements, as well as defect scattering matrix elements are calculated using the fifth-nearest neighbor tight-binding model, as described in ref 37 on a  $360 \times 360$  grid of  $\mathbf{k}$ -points. The phonon dispersions and the phonon eigenvectors needed for the calculation of the  $e-ph$  matrix elements are derived using density-functional perturbation theory as implemented in the Quantum ESPRESSO<sup>38</sup> DFT package within local density approximation (LDA).<sup>35</sup> Norm-conserving pseudopotentials with a 55 Ry cutoff for the wave function are used, and the BZ is sampled by  $32 \times 32$   $\mathbf{k}$ -points in the calculation of the electronic states. The phonons are calculated on a  $8 \times 8$   $\mathbf{q}$ -point grid and, for the Raman intensity, they are interpolated into a  $120 \times 120$  grid. The  $\delta$  functions in eq 6 are broadened into Lorentzians with FWHM  $\sim 8 \text{ cm}^{-1}$  to compensate for the finite computational grid.<sup>37</sup> As defect, a weakened nearest-neighbor tight-binding hopping element is used, following ref 37, with the perturbation being  $\delta t$ . The amount of defects is characterized by a parameter

$\alpha_{\text{hopp}}$  describing both defect density and magnitude of defect perturbation,  $\alpha_{\text{hopp}} = \delta t n_D = 6.4 \times 10^{13} \text{ eV}^2 \text{ cm}^{-2}$ , as for ref 37. A value of 68 meV for the doping-independent part of the broadening, consisting of contributions due to  $e-ph$  and  $e-def$  scattering, can well reproduce the experimentally observed trend of a decrease of the areas of the D and 2D peaks, as shown in Figure 9. The  $e-ph$  contribution in the present samples is estimated from our experimental data to be  $\sim 31$  meV, following the procedure of ref 21. This implies that the defect part is  $\approx 40$  meV. To the best of our knowledge, this value has never been determined experimentally. The defect-related

broadening depends on the number of defects and also the type of defects. Note that there is a difference of a factor of 4 between the definition of  $\gamma$  of ref 37 and that used here, chosen to be consistent with the notation in ref 21.

Figure 9 plots the doping dependence of  $A(D)/A(G)$ ,  $A(2D)/A(G)$  and  $A(D)/A(2D)$  for sample A, normalized with respect to their values at low doping. The area of the G peak is constant with doping for  $\hbar\omega_L/2 > E_F$ ,<sup>21</sup> therefore, the dependencies of  $A(D)/A(G)$  and  $A(2D)/A(G)$  with doping are representative of the  $A(D)$  and  $A(2D)$  trend, respectively. The agreement between experiment and theory is remarkable, given the simple description of the doping in the simulation. The increase of the total broadening due to  $\gamma_{ee}$ , as described in ref 21, can well reproduce the trend of the decrease of the peaks' areas, strongly indicating that  $e-e$  correlation is likely to be the most relevant cause for the observed experimental trends. As the concentration of charge carriers is increased, some of the other ingredients of the model, such as electron and phonon dispersions, might change. With doping, the graphene lattice parameter is expected to change, leading to a shift in the adiabatic energy of the phonon modes.<sup>20,29</sup> This changes the position of the Raman peaks, but should not significantly alter their intensity. The  $e-ph$  matrix elements are also expected to decrease slightly

with increasing doping,  $\sim 15\%$  for the doping range probed in this experiment,<sup>39</sup> thus representing a correction to our analysis. More detailed studies are required to fully understand the interplay between the different effects.

## CONCLUSIONS

We studied the dependence of Raman spectrum of defected graphene on the charge carrier concentration, by combining polymer electrolyte gating with *in situ* Hall-effect and Raman measurements. For a given number of defects, the intensities and areas of the D and D' peaks decrease with increasing doping. Considering all Raman scattering processes, we assigned the doping-induced intensity variation to the increased total scattering rate of the photoexcited electrons and holes, resulting from the doping-dependent strength of electron-electron scattering. This analysis paves the way for the experimental evaluation of the different sources of broadening in the electronic states in graphene and a better understanding of the role and type of defects on its physical properties. This highlights the importance of taking into account the doping level when determining the amount and the type of defects from the intensity of the D peak. We thus presented general relations between D peak intensity and defects valid for any doping level.

## METHODS

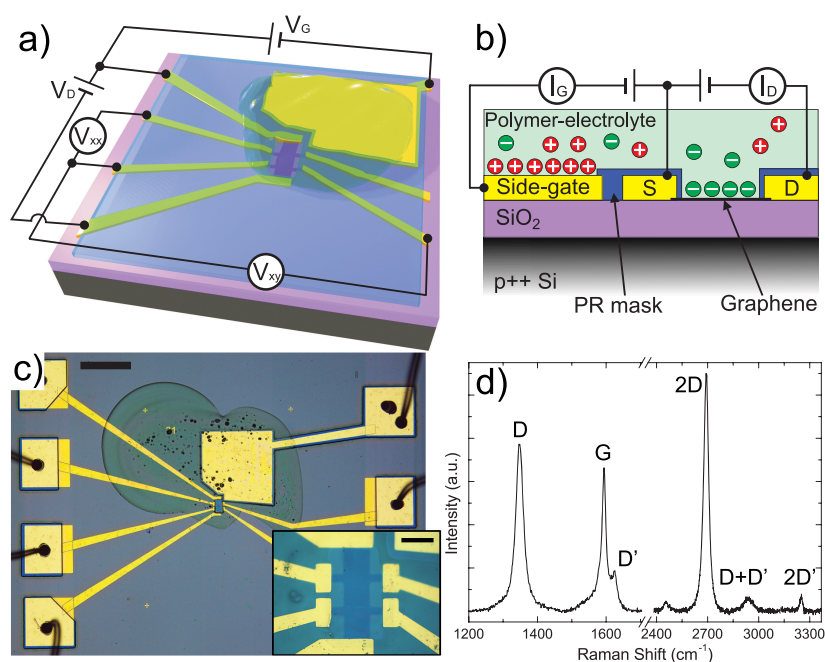
The defected samples are prepared as follows. A SLG film is grown by chemical vapor deposition (CVD) on a 25  $\mu\text{m}$  copper foil.<sup>27,30</sup> It is then transferred on a Si + 300 nm  $\text{SiO}_2$  substrate as described in refs 27 and 40. The Hall bar geometry is defined by creating a photoresist (PR) mask by photolithography and removing the uncovered portion of the film by  $\text{O}_2/\text{Ar}$  reactive ion etching. A further PR layer is then spun on the sample and windows are opened by photolithography only on the graphene channel and side-gate area (see Figure 10). The PR mask is then hard-baked at 145  $^\circ\text{C}$  for 5 min in order to improve its chemical stability.<sup>41</sup> The devices are then exposed to a mild  $\text{O}_2$  inductively coupled plasma to introduce defects in the graphene channel.<sup>42</sup> The process is carried out at a pressure  $\sim 150$  mTorr, with a power of 15 W and for a few seconds (typically between 2 and 10s) depending on the desired defect concentration.

$E_F$  is moved by applying a droplet of polymer electrolyte, consisting of  $\text{LiClO}_4$  and poly(ethylene oxide) in the weight ratio 0.12:1, over both the device channel and the side-gate.<sup>14</sup> The working principle of this gating technique is shown schematically in Figure 10a,b. When a potential is applied between the side-gate and the graphene channel, free ions migrate and accumulate at the surface of the electrodes to form an electric double layer (EDL). Because of the large interfacial capacitance of the EDL, the compensation of these charges shifts  $E_F$ , much more than what can be achieved with standard dielectric gates.<sup>14,20,28</sup> In the case of standard dielectric-gating, the applied gate potential uniformly drops across the gate dielectric, and it is therefore possible to directly estimate the induced charge through a capacitor model.<sup>20,44</sup> In a polymer-electrolyte gated field-effect transistor, the applied gate potential drops across two nanocapacitors in series (one at the side-gate/electrolyte interface, the other at the electrolyte/channel interface), separated by an ionic conductive medium. In order

to maximize the voltage drop across the channel/electrolyte interface, the side-gate/electrolyte capacitance must be the dominant one, hence the area of the gate electrode is significantly larger than that of the graphene channel, and a PR mask is fabricated on top of the device in order to minimize the direct contact area between metal electrodes and polymer electrolyte, thus reducing parasitic capacitance. Despite this, a voltage drop at the gate electrode cannot be excluded, and might lead to errors when correlating the applied gate voltage to the amount of induced charges in the graphene channel and  $E_F$ .<sup>45</sup> To avoid possible systematic errors,  $E_F$  in the graphene channel is directly evaluated by Hall-effect measurements. A 1  $\mu\text{A}$  direct current (DC) is applied between source and drain leads using a Keithley 2410 source-measurement unit, while longitudinal and transverse voltages ( $V_{xx}$  and  $V_{xy}$ ) are measured by means of a Keithley 2182A nanovoltmeter. The perpendicular magnetic field is applied using a permanent magnet with a surface field of 0.37T, as measured using a calibrated Gauss-meter. Another Keithley 2410 is used to apply the gate voltage,  $V_g$ . The overall performance of polymer-electrolyte gating is limited by the electrochemical stability of the polymer-electrolyte. When a  $V_g$  higher than the electrochemical stability window of the polymer electrolyte is applied, electrochemical reactions, such as hydrolysis of residual water in the electrolyte,<sup>46,47</sup> can occur and permanently modify the graphene electrode, thus changing the total amount of defects. In order to avoid this, the maximum applied  $V_g \sim \pm 2$  V and the gate leakage current, a good indication of possible electrochemical reactions,<sup>47,48</sup> is monitored and kept at  $\sim 10^{-10}$  A. The absence of permanent modifications is confirmed by the repeatability of the Raman measurements through several gate voltage sweeps.

Raman measurements are carried out at room temperature in a Renishaw InVia microspectrometer equipped with a 100 $\times$  objective (numerical aperture 0.9). The spot size is  $\sim 1$   $\mu\text{m}$ , and the incident power is kept well below 1 mW in order to avoid





**Figure 10.** (a) Scheme of the polymer electrolyte-gated graphene transistor and electrical configuration used for the measurements. (b) Polymer electrolyte gating process. When biasing graphene with respect to the side-gate, Li<sup>+</sup> (red) and ClO<sub>4</sub><sup>-</sup> (green) ions migrate to form electric double layers near each electrode.<sup>14,43</sup> (c) Optical micrograph of the device used in our experiments, scale bar 300 μm. Inset: graphene channel; scale bar 30 μm. (d) Raman spectrum of defected SLG.

heating effects. The excitation wavelengths are 514 nm for an Ar<sup>+</sup> laser and 633 nm from a He–Ne laser, chosen because these are the most commonly used for Raman characterization of graphene.<sup>2,3,11</sup>

**Conflict of Interest:** The authors declare no competing financial interest.

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